

Transport Calculations with TranSiesta

Pablo Ordejón

Instituto de Ciencia de Materiales de Barcelona - CSIC, Spain



M. Brandbyge, K. Stokbro and J. Taylor

Mikroelektronik Centret - Technical Univ. Denmark



Jose L. Mozos and Frederico D. Novaes

Instituto de Ciencia de Materiales de Barcelona - CSIC, Spain



The SIESTA team

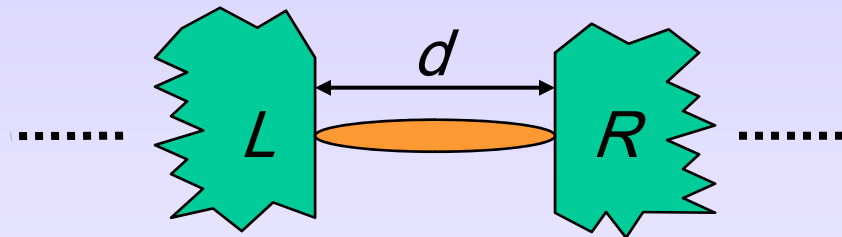
E. Anglada, E. Artacho, A. García, J. Gale. J. Junquera, D. Sánchez-Portal,
J. M. Soler,



Outline

1. Electronic transport in the nanoscale: Basic theory
2. Modeling: the challenges and our approach:
 - The problem at equilibrium (zero voltage)
 - Non-equilibrium (finite voltage)
3. Practicalities

1. Electronic Transport in the Nanoscale: Basic Theory

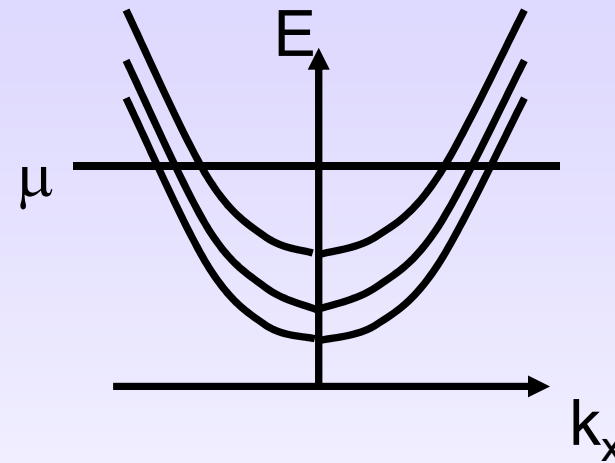
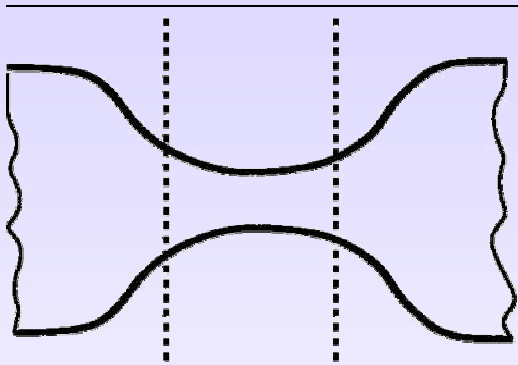


- Scattering in nano-scale systems:
 - electron-electron interactions
 - phonons
 - impurities, defects
 - elastic scattering by the potential of the contact
- Semiclassical theory breaks down – QM solution needed
- Landauer formulation: Conductance as transmission probability

$$\begin{aligned} d &\ll L_m \\ d &\approx \lambda \end{aligned}$$

S. Datta, *Electronic transport in mesoscopic systems* (Cambridge)

Narrow constriction (meso-nanoscopic)

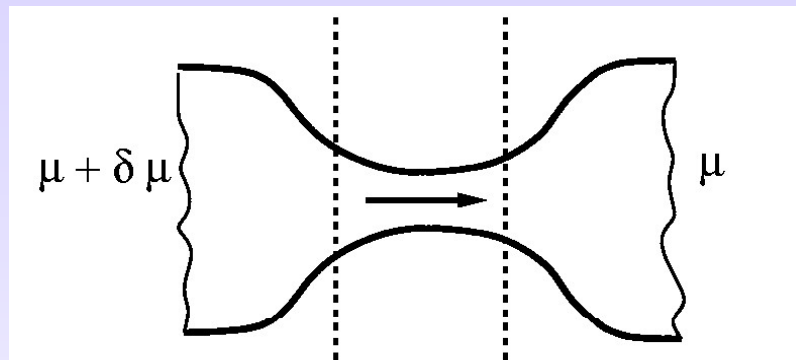


Transversal confinement $\Rightarrow$ QUANTIZATION

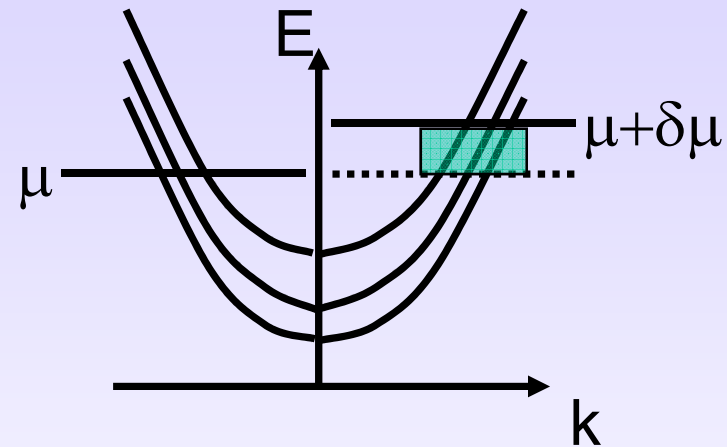
$$\psi(x, y) = \psi_m(y) e^{ik_x x}$$

$$E_m(k_x) = \varepsilon_m + \frac{\hbar^2 k_x^2}{2m}$$

Landauer formulation - no scattering



$$\delta\mu = -eV$$



$$I = -2e \sum_i^{\text{bands}} \int_{\text{BZ}} dk f(k) v_g(k, i) \approx \frac{2e^2}{h} V N_{\text{bands}}(E_F)$$

$$G = \frac{I}{V} = N_{\text{bands}} \frac{2e^2}{h} = N_{\text{bands}} G_0$$

QUANTUM OF CONDUCTANCE

$$G_0 = \frac{2e^2}{h} \Rightarrow 12.9 \text{ k}\Omega$$

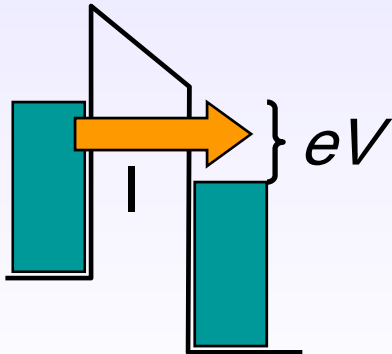
Landauer formulation - scattering

- Transmission probability of an incoming electron at energy ε :

$$T(\varepsilon) = \text{Tr}[t^\dagger t](\varepsilon)$$

transmission matrix: $\psi_{out} = t \psi_{in}$

- Current:

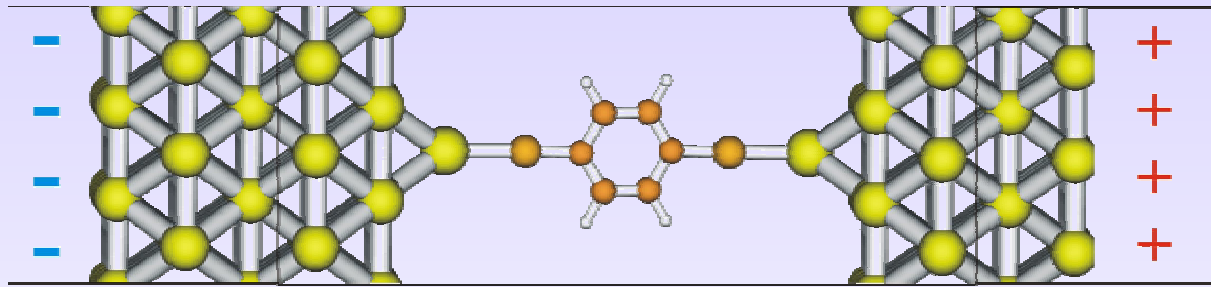


$$I = \frac{2e^2}{h} \int d\varepsilon (f_L(\varepsilon) - f_R(\varepsilon)) T(\varepsilon)$$

- Perfect conductance (one channel): $T=1$

$$G_0 = \frac{2e^2}{h}$$

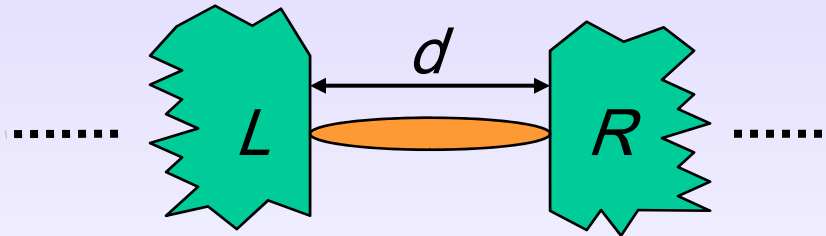
2. Modeling: The Challenges and Our Approach



- Model the molecule-electrode system from *first principles*:
No parameters fitted to the particular system → DFT
- Model a molecule coupled to bulk (infinite) electrodes
- Electrons out of equilibrium (do not follow the thermal Fermi occupation)
- Include finite bias voltage/current and determine the potential profile
- Calculate the conductance (quantum transmission through the molecule)
- Determine geometry: Relax the atomic positions to an energy minimum

Restriction: Ballistic conduction

- **Ballistic conduction**: consider only the scattering of the incoming electrons by the potential created by the contact



- **Two terminal devices** (three terminals in progress)
- Effects not described: inelastic scattering
 - electron-electron interactions (Coulomb blockade)
 - phonons (current-induced phonon excitations)

First Principles: DFT

- Many interacting-electrons problem mapped to a one particle problem in an external effective potential (Hohenberg-Kohn, Kohn-Sham)

$$V_{eff} = V_{ext} + V_{ps} + V_H + V_{XC}$$

- Charge density as basic variable:

$$V_{eff} = V_{eff}[\rho(r)]$$

- Self-consistency: $V_{eff} \leftrightarrow \rho$
- Ground state theory: V_{XC}



SIESTA

<http://www.uam.es/siesta>

Soler, Artacho, Gale, García, Junquera, Ordejón and Sánchez-Portal
J. Phys.: Cond. Matt. 14, 2745 (2002)

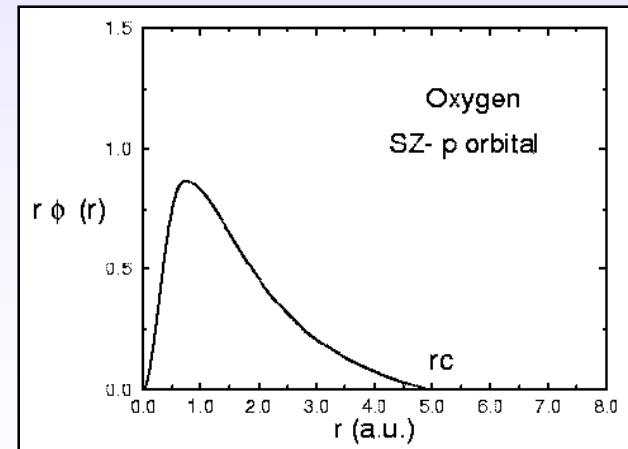
- Self-consistent DFT code (LDA, GGA, LSD)
- Pseudopotentials (Kleinman-Bylander)
- LCAO approximation:

Basis set:

Confined Numerical Atomic Orbitals

$$\varphi_{\mu}(\mathbf{r})$$

Sankey's "fireballs"



- Order-N methodology (in the calculation and the solution of the DFT Hamiltonian)

TRANSIESTA

Implementation of non-equilibrium electronic transport in SIESTA

- Atomistic description (both contact and electrodes)
- Infinite electrodes
- Electrons out of equilibrium
- Include finite bias and determine the potential profile
- Calculates the conductance (both linear and non-linear)
- Forces and geometry

Brandbyge, Mozos, Ordejón, Taylor and Stokbro
Phys. Rev. B. **65**, 165401 (2002)

Mozos, Ordejón, Brandbyge, Taylor and Stokbro
Nanotechnology **13**, 346 (2002)

The problem at Equilibrium (Zero Bias)

Challenge:

Coupling the finite contact to infinite electrodes

Solution:

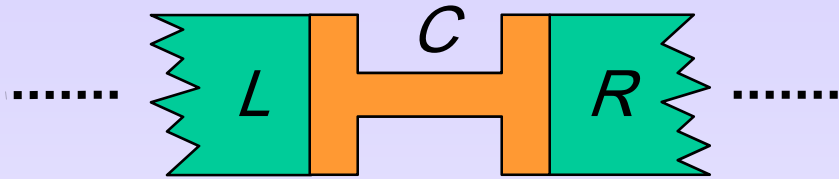
Green's Functions

$$G(z) = (z - H)^{-1}$$

$$\rho(\varepsilon) = -\frac{1}{\pi} \text{Im} G(\varepsilon + i\delta)$$

$$\rho(\mathbf{r}) = \sum_{\mu\nu} \underbrace{\left[\int_{-\infty}^{\infty} d\varepsilon \rho_{\mu\nu}(\varepsilon) n_F(\varepsilon - \mu_F) \right]}_{D_{\mu\nu}} \varphi_{\mu}(\mathbf{r}) \varphi_{\nu}(\mathbf{r})$$

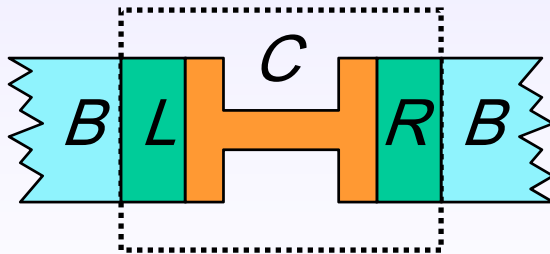
Setup (zero bias)



$$G(\varepsilon) = (\varepsilon - H)^{-1} = \begin{bmatrix} \dots & -V_L^+ & & & \\ -V_L^- & \varepsilon - H_L & -V_{LC} & & \\ & -V_{CL} & \varepsilon - H_C & -V_{CR} & \\ & & -V_{RC} & \varepsilon - H_R & -V_R^+ \\ & & & -V_R^- & \dots \end{bmatrix}^{-1}$$

Contact:

- Contains the molecule, and part of the Right and Left electrodes
- Sufficiently large to include the screening



Solution in finite system:

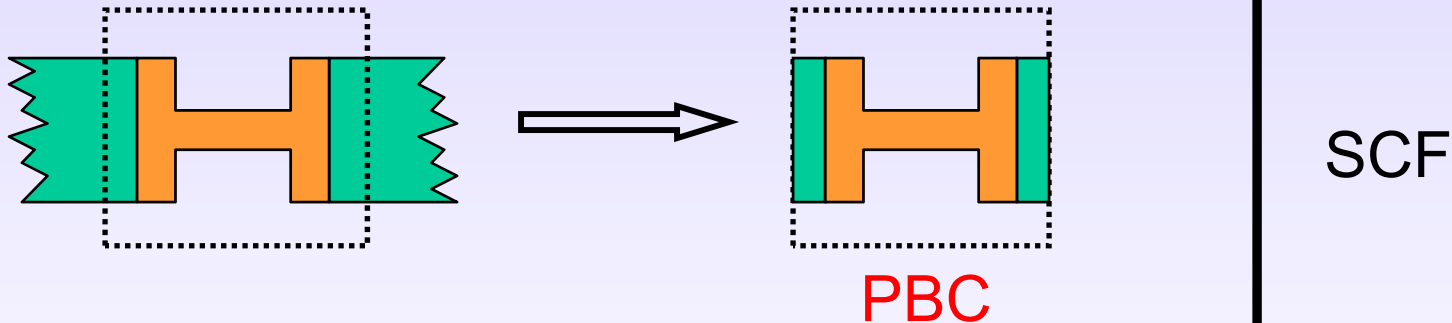
$$G(\varepsilon) = \begin{bmatrix} \varepsilon - H_L - \Sigma_L & -V_{LC} & \\ -V_{CL} & \varepsilon - H_C & -V_{CR} \\ & -V_{RC} & \varepsilon - H_R - \Sigma_R \end{bmatrix}^{-1}$$

$\Sigma(\varepsilon)$ = Selfenergies. Can be obtained from the bulk Greens functions

Lopez-Sancho et al. J. Phys. F 14, 1205 (1984)

Calculations (zero bias):

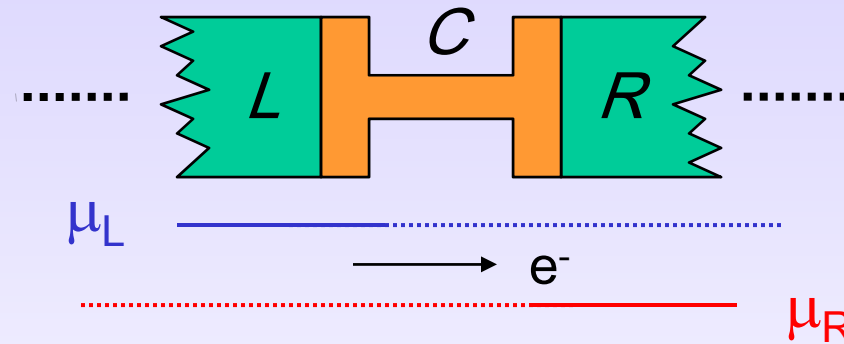
- Bulk Greens functions and self-energies (unit cell calculation)
- Hamiltonian of the Contact region:



- Solution of GF's equations $\Rightarrow \rho(r)$
- Landauer-Büttiker: transmission probability: $T(\varepsilon) = Tr[t^\dagger t](\varepsilon)$

$$t(\varepsilon) = \left(Im[\Sigma_R(\varepsilon)] \right)^{1/2} G(\varepsilon) \left(Im[\Sigma_L(\varepsilon)] \right)^{1/2}$$

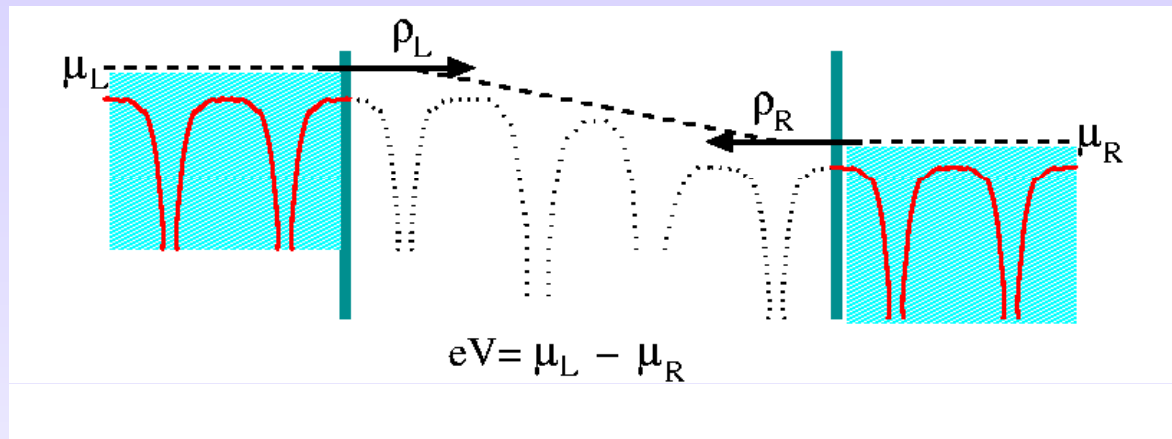
The problem at Non-Equilibrium (Finite Bias)



2 additional problems:

- Non-equilibrium situation:
 - current flow
 - two different chemical potentials
- Electrostatic potential with boundary conditions

Non-equilibrium formulation:



- Scattering states (from the left)
Lippmann-Schwinger Eqs.:

$$\psi_l = \psi_l^o + G(\varepsilon_o) V_L \psi_l^o$$

- Non-equilibrium Density Matrix:

$$D_{\mu\nu} = \int_{-\infty}^{\infty} d\varepsilon \left[\rho_{\mu\nu}^L(\varepsilon) n_F(\varepsilon - \mu_L) + \rho_{\mu\nu}^R(\varepsilon) n_F(\varepsilon - \mu_R) \right]$$

$$\rho_{\mu\nu}^L(\varepsilon) = \frac{1}{\pi} \left(\boxed{G}(\varepsilon + i\delta) \text{Im} \left[\boxed{\Sigma_L}(\varepsilon + i\delta) \right] \boxed{G}^\dagger(\varepsilon + i\delta) \right)_{\mu\nu}$$

Electrostatic Potential

Given $\rho(r)$, $V_H(r)$ is determined except up to a linear term:

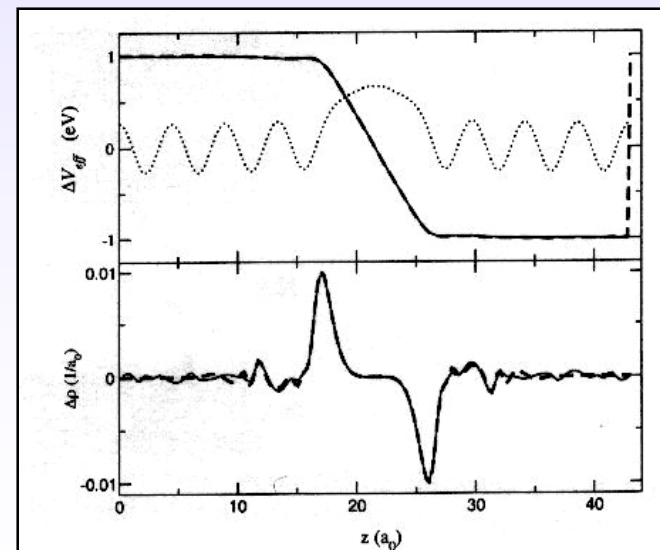
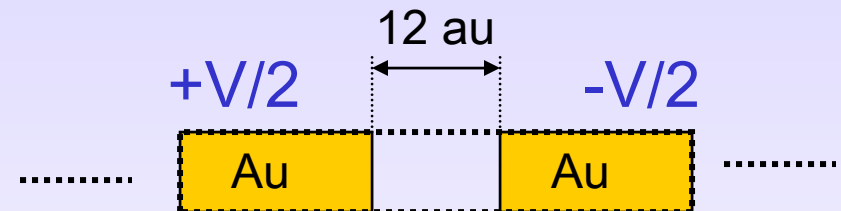
$$V_H(\mathbf{r}) = \phi(\mathbf{r}) + \mathbf{a} \cdot \mathbf{r} + b$$

$\phi(r)$: particular solution of Poisson's equation

a and b : determined imposing BC: the shift V between electrodes

- $\phi(r)$ computed using FFT's

- Linear term: $-\frac{V}{L} \left(z - \frac{L}{2} \right)$

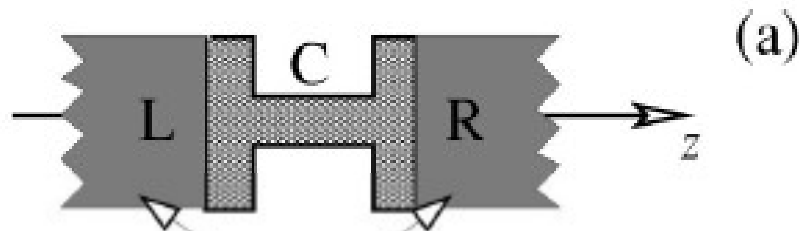


3. TranSIESTA Practicalities

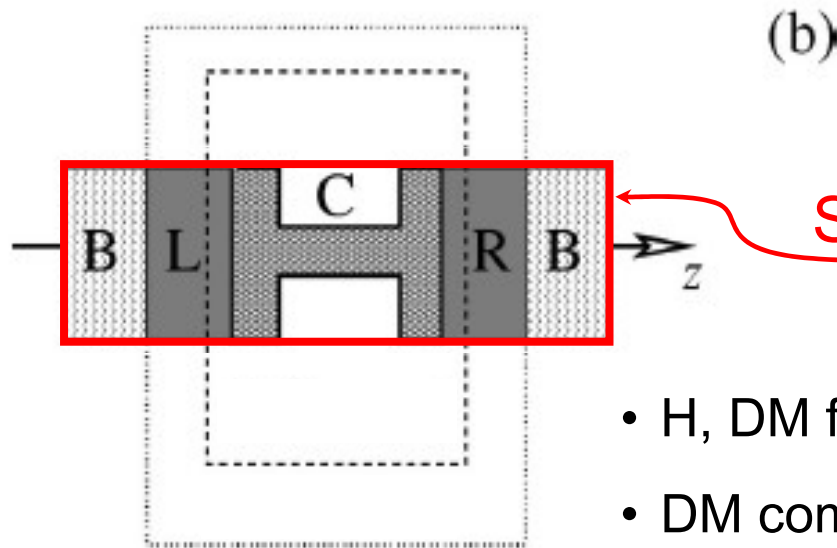
3 Step process:

1. SIESTA calculation of the bulk electrodes, to get H , ρ , and Self-energies
2. SIESTA calculation for the open system
 - reads the electrode data
 - builds H from ρ
 - solves the open problem using Green's Functions (TranSIESTA)
 - builds new ρ
3. Postprocessing: compute $T(E)$, I , ...





Semi-infinite bulk

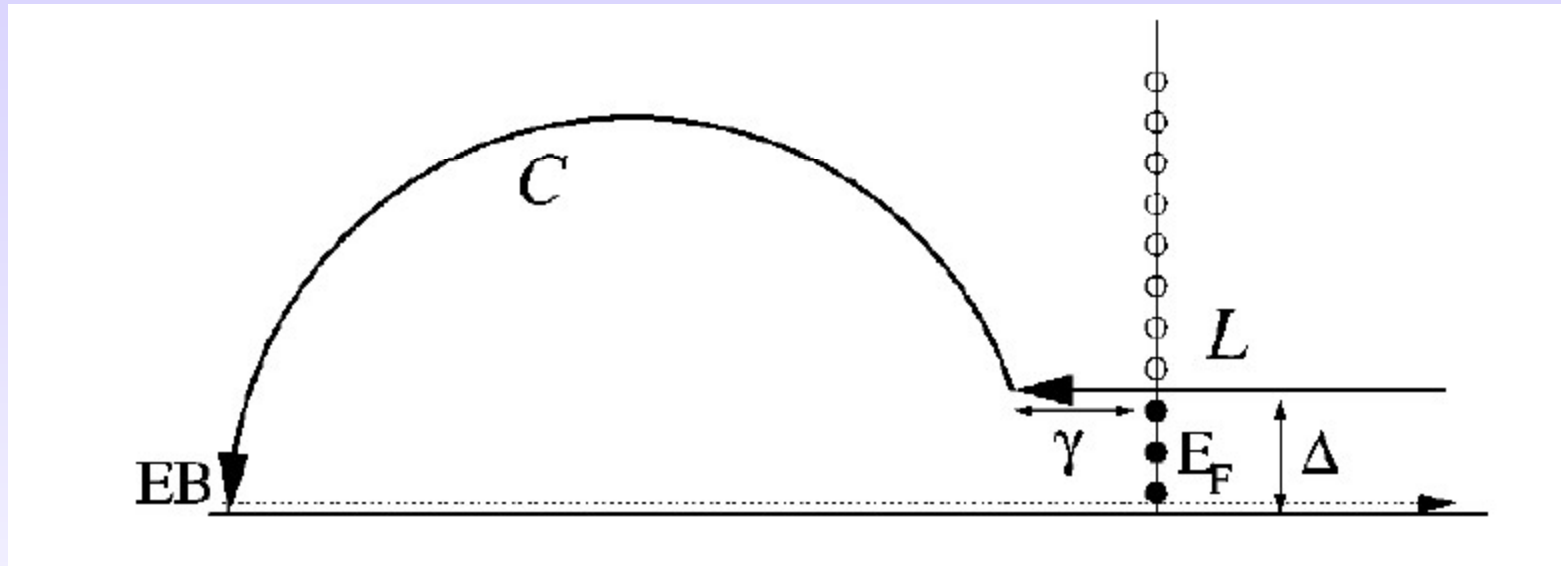


Supercell - PBC

$$G(\varepsilon) = \begin{bmatrix} \varepsilon - H_L - \Sigma_L & -V_{LC} & 0 \\ -V_{CL} & \varepsilon - H_C & -V_{CR} \\ 0 & -V_{RC} & \varepsilon - H_R - \Sigma_R \end{bmatrix}^{-1}$$

- H , DM fixed to bulk in L and R
- DM computed in C from Green's functions
- H_C , V_{LC} and V_{CR} computed in a supercell approach (with potential ramp)
- B (buffer) does not enter directly in the calculation (only in the SC calc. for V_{Hartree})

Contour integration



$$\begin{aligned} \mathbf{D} &= -\frac{1}{\pi} \int_{-\infty}^{\infty} d\varepsilon \operatorname{Im}[\mathbf{G}(\varepsilon + i\delta)] n_F(\varepsilon - \mu) \\ &= -\frac{1}{\pi} \operatorname{Im} \left[\int_{-\infty}^{\infty} d\varepsilon \mathbf{G}(\varepsilon + i\delta) n_F(\varepsilon - \mu) \right]. \end{aligned}$$

Contour Integration

$$\mathbf{D}_{\mu\nu} = \mathbf{D}_{\mu\nu}^L + \Delta_{\mu\nu}^R, \quad (28)$$

$$\mathbf{D}_{\mu\nu}^L = -\frac{1}{\pi} \text{Im} \left[\int_{EB}^{\infty} d\varepsilon \mathbf{G}(\varepsilon + i\delta) n_F(\varepsilon - \mu_L) \right], \quad (29)$$

$$\Delta_{\mu\nu}^R = \int_{-\infty}^{\infty} d\varepsilon \rho_{\mu\nu}^R(\varepsilon) [n_F(\varepsilon - \mu_R) - n_F(\varepsilon - \mu_L)], \quad (30)$$

or equivalently

$$\mathbf{D}_{\mu\nu} = \mathbf{D}_{\mu\nu}^R + \Delta_{\mu\nu}^L, \quad (31)$$

$$\mathbf{D}_{\mu\nu}^R = -\frac{1}{\pi} \text{Im} \left[\int_{EB}^{\infty} d\varepsilon \mathbf{G}(\varepsilon + i\delta) n_F(\varepsilon - \mu_R) \right], \quad (32)$$

$$\Delta_{\mu,\nu}^L = \int_{-\infty}^{\infty} d\varepsilon \rho_{\mu\nu}^L(\varepsilon) [n_F(\varepsilon - \mu_L) - n_F(\varepsilon - \mu_R)]. \quad (33)$$

$$\mathbf{D}_{\mu\nu} = w_{\mu\nu} (\mathbf{D}_{\mu\nu}^L + \Delta_{\mu\nu}^R) + (1 - w_{\mu\nu}) (\mathbf{D}_{\mu\nu}^R + \Delta_{\mu\nu}^L), \quad (36)$$

$$w_{\mu\nu} = \frac{(\Delta_{\mu\nu}^L)^2}{(\Delta_{\mu\nu}^L)^2 + (\Delta_{\mu\nu}^R)^2}. \quad (37)$$

SolutionMethod Transiesta

```
#          GENGf  OPTIONS
TS.ComplexContour.Emin    -3.0 Ry
TS.ComplexContour.NPoles      6
TS.ComplexContour.NCircle    20
TS.ComplexContour.NLine      3
TS.RealContour.Emin        -3.0 Ry
TS.RealContour.Emax         2.d0 Ry
TS.TBT.Npoints              100
```

```
# TS OPTIONS
TS.Voltage 1.000000 eV
TS.UseBulkInElectrodes .True.
TS.BufferAtomsLeft  0
TS.BufferAtomsRight 0
```

```
# TBT OPTIONS
TS.TBT.Emin -5.5 eV
TS.TBT.Emax +0.5 eV
TS.TBT.NPoints 100
TS.TBT.NEigen 3
```

