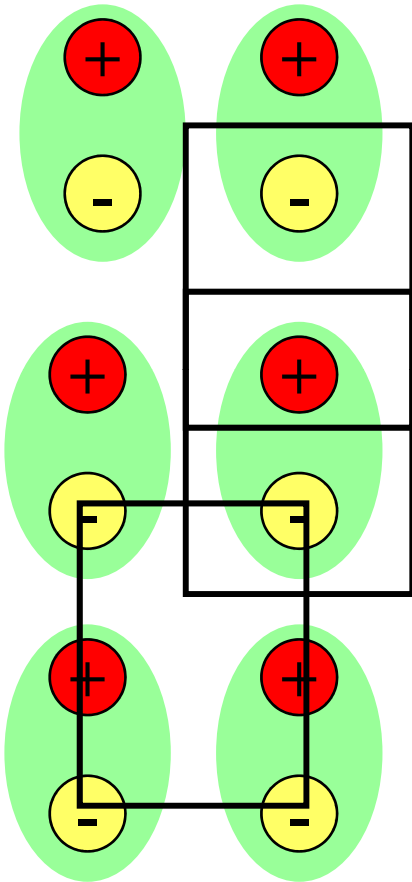


Electric polarization



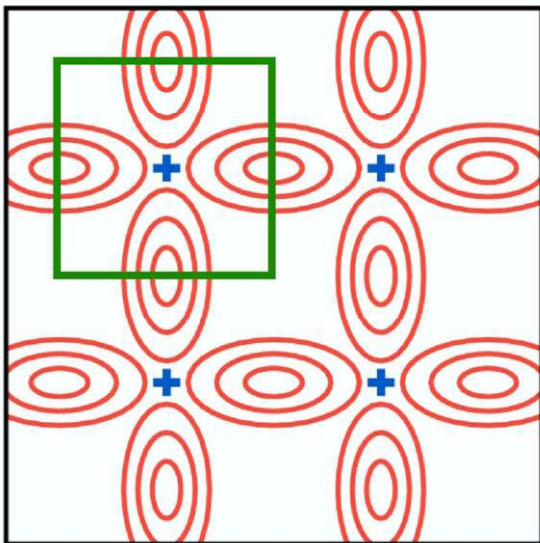
How to define the electric polarization in a solid?

Compose of molecules: $\vec{P} = \vec{d}_{mol} / V_{mol}$

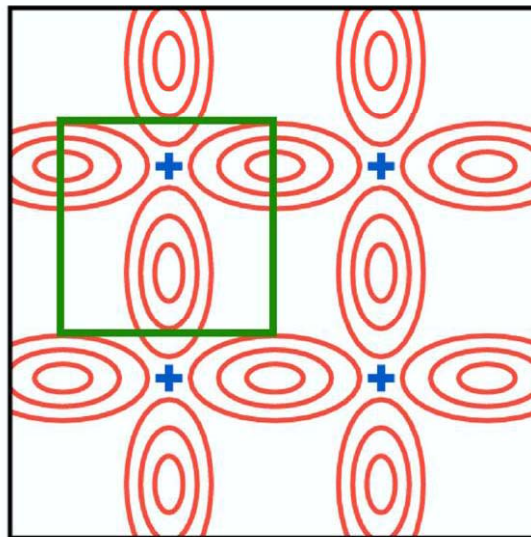
Electric polarization cannot uniquely be defined from the charge distribution.

For a covalent solid the problem is even more clear...

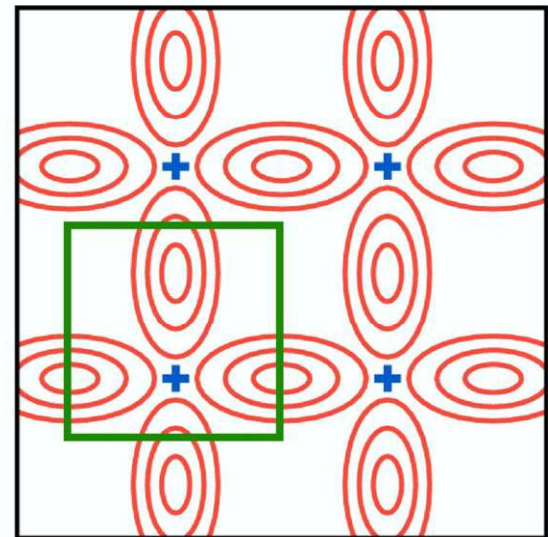
$d = 0$



$d > 0$



$d < 0$



Correct way: use currents

Use the current instead of density

$$\frac{d\vec{P}_e}{dt} = \frac{1}{V_{cell}} \int d\vec{r}^3 j(\vec{r})$$

R. Resta 1992: changes in the electric polarization can be unambiguously defined and calculated

$$\Delta\vec{P}_e = \int dt \frac{d\vec{P}}{dt}$$

How to calculate polarization changes

King-Smith and Vanderbilt, 1993:

$$\Delta \mathbf{P} = \mathbf{P}(t_2) - \mathbf{P}(t_1) \quad \text{where}$$

$$\mathbf{P} = \frac{ie}{(2\pi)^3} \sum_n \int_{\text{BZ}} d^3k \langle u_{nk} | \nabla_{\mathbf{k}} | u_{nk} \rangle$$

where $\psi_{nk}(\mathbf{r}) = e^{i\mathbf{k} \cdot \mathbf{r}} u_{nk}(\mathbf{r})$

↑
“Berry phase”

Ionic contribution

Of course, we have to add the contribution due the core of atoms or to the ions (i.e., core + core electrons):

$$\Delta \vec{P}_{tot} = \Delta \vec{P}_e + \sum_i Z^i_{val} \Delta \vec{R}_i$$

Equivalent formulation in terms of Wannier functions

$$\vec{P}_e = -\frac{2e}{V_{cell}} \sum_{i=1}^{N/2} \langle W_i | \vec{r} | W_i \rangle$$

The centers of the Wanniers are defined modulo \vec{R}

$$\vec{r}_i \rightarrow \vec{r}_i + \vec{R} \quad \text{with} \quad \vec{R} = n_1 \vec{a}_1 + n_2 \vec{a}_2 + n_3 \vec{a}_3$$

and, therefore, the polarization is defined modulo a “quantum” of polarization.

$$\vec{P}_e \rightarrow \vec{P}_e \pm \frac{2e\vec{R}}{V_{cell}}$$

Why is interesting to calculate ΔP ?

***You can study ferroelectricity**

***You can calculate the dynamical or effective Born charges, Z**

$$Z_{j\alpha\beta}^* = \frac{dP_\alpha}{dR_{j\beta}} \simeq \frac{\Delta P_\alpha}{\Delta R_{j\beta}}$$

***With Z you can calculate the IR absorption**

***With Z you can calculate the long-range part of the dynamical matrix and LO-TO splitting**

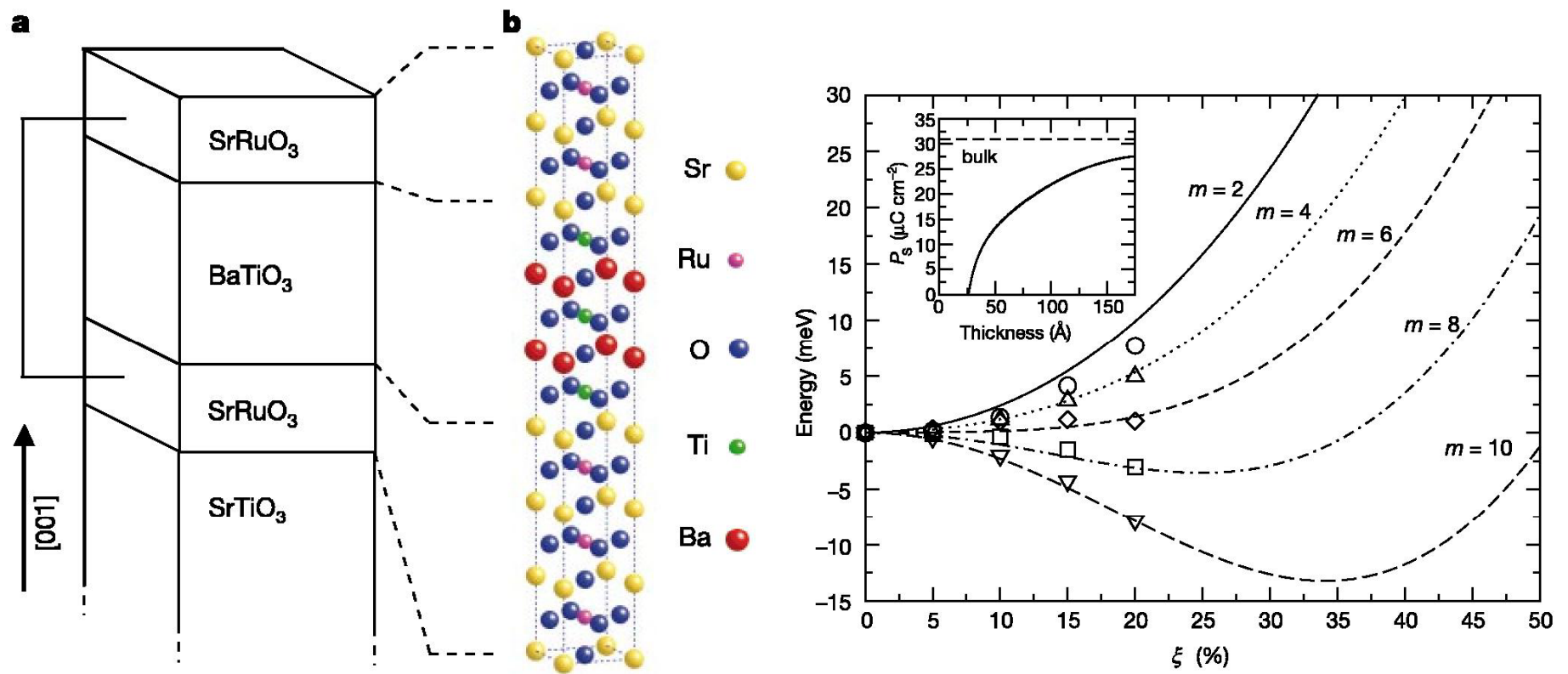
$$D_{mn}^{\mu\nu}(\mathbf{q}) = D_{mn}^{(0)\mu\nu}(\mathbf{q}) + \frac{4\pi e^2}{\Omega} \frac{(Z_m^* \cdot \hat{\mathbf{q}})_\mu (Z_n^* \cdot \hat{\mathbf{q}})_\nu}{\epsilon_\infty(\mathbf{q})},$$

Dynamical charges in BaTiO₃

	BaTiO ₃		
	SIESTA	PW ²³	PW ²⁷
Z_A^*	+2.72	+2.74	+2.75
Z_{Ti}^*	+7.60	+7.32	+7.16
$Z_{O_{II}}^*$	-2.18	-2.14	-2.11
$Z_{O_I}^*$	-5.96	-5.78	-5.69

Critical thickness of ferroelectric thin films

J. Junquera and P. Ghosez, Nature 422, 506 (2003)



SIESTA INPUT

PolarizationGrids:

$$P_{e,\parallel} = \frac{ifq_e}{8\pi^3} \int_A d\mathbf{k}_\perp \sum_{n=1}^M \int_0^{|G_\parallel|} dk_\parallel \langle u_{\mathbf{k}n} | \frac{\delta}{\delta k_\parallel} | u_{\mathbf{k}n} \rangle$$

```
%block PolarizationGrids
```

```
10  4  4
```

```
4  10  4
```

```
4  4  10
```

```
%endblock PolarizationGrids
```

SIESTA INPUT

BornCharge .true. (together with MD.TypeOfRun = FC)

The Born effective charge matrix is then written to the file SystemLabel.BC

**We aware of the ambiguous definition of the polarization
in the numerical calculation.....**

